Snowmass2013

Community Planning Meeting aka CPM

High Energy Frontier

Michael Peskin (SLAC)
Chip Brock (MSU)

TOC:

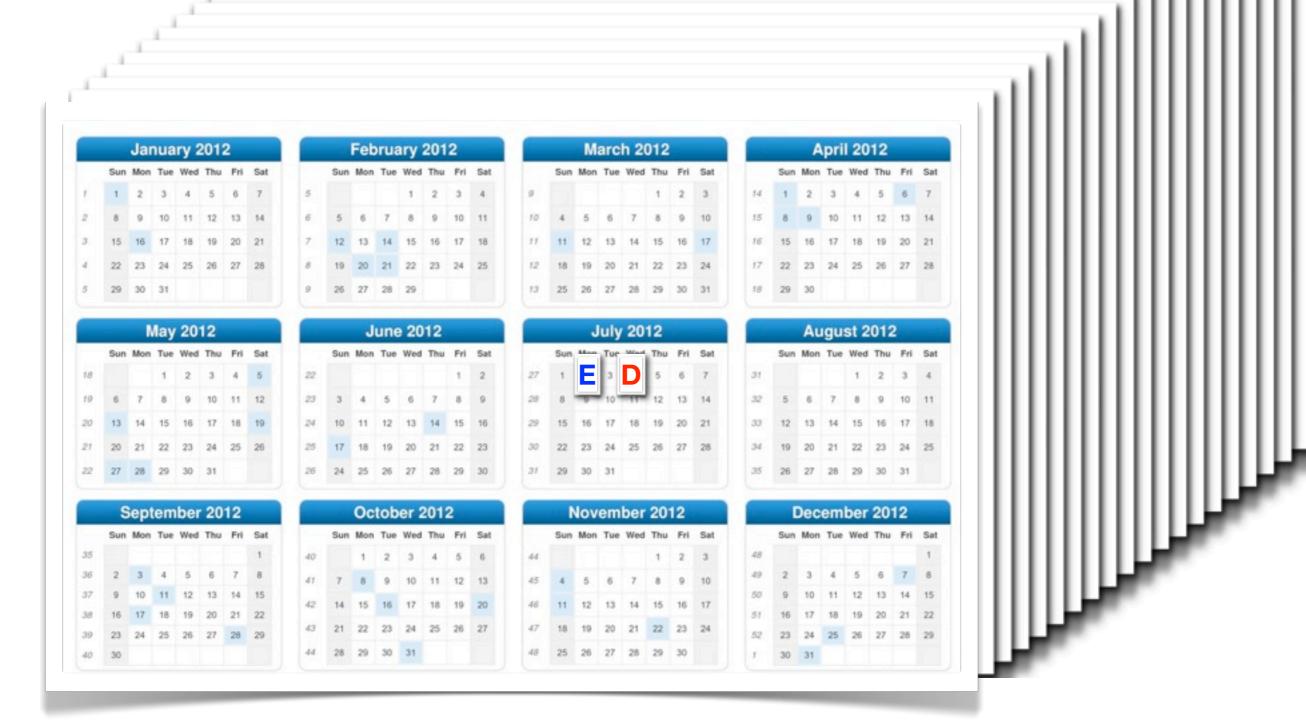
- 1. High Energy Frontier, introduction
- 2. High Energy Frontier Study, practicalities
- 3. What's next for the High Energy Frontier
- 4. This Workshop

Snowmass is about: following the physics

and the route has just gotten a little clearer.

From 1967 to July

It's been a long 45 years







W LETTERS

20 NOVEMBER 1967

VOLUME 19, NUMBER 21

PHYSICAL REVIEW LETTERS

you never get tired of.

11 In obtaining the expression (11) the mass difference between the charged and neutral has been ignored.

¹²M. Ademollo and R. Gatto, Nuovo Cimento 44A, 282 (1966); see also J. Pasupathy and R. E. Marshak, Phys. Rev. Letters <u>17</u>, 888 (1966).

¹³The predicted ratio [eq. (12)] from the current alge-

bra is slightly larger than that (0.23%) obtained from the ρ -dominance model of Ref. 2. This seems to be true also in the other case of the ratio $\Gamma(\eta \to \pi^*\pi^*\gamma)/\Gamma(\gamma)$ calculated in Refs. 12 and 14. ¹⁴L. M. Brown and P. Singer, Phys. Rev. Letters 8,

A MODEL OF LEPTONS

Steven Weinberg†

Laboratory for Nuclear Science and Physics Department, Massachusetts Institute of Technology, Cambridge, Massachusetts (Received 17 October 1967)

Leptons interact only with photons, and with the intermediate bosons that presumably mediate weak interactions. What could be more natural than to unite1 these spin-one bosons into a multiplet of gauge fields? Standing in the way of this synthesis are the obvious differences in the masses of the photon and intermediate meson, and in their couplings. We might hope to understand these differences by imagining that the symmetries relating the weak and electromagnetic interactions are exact symmetries of the Lagrangian but are broken by the vacuum. However, this raises the specter of unwanted massless Goldstone bosons. This note will describe a model in which the symmetry between the electromagnetic and weak interactions is spontaneously broken, but in which the Goldstone bosons are avoided by introducing the photon and the intermediateboson fields as gauge fields.3 The model may

We will restrict our attention to symmetry groups that connect the observed electron-type leptons only with each other, i.e., not with muon-type leptons or other unobserved leptons or hadrons. The symmetries then act on a left-

$$L = \left[\frac{1}{2}(1 + \gamma_5)\right]\begin{pmatrix} e \\ e \end{pmatrix}$$
(1)

and on a right-handed singlet

$$R = \left[\frac{1}{2}(1-\gamma_5)\right]e$$
. (2)

The largest group that leaves invariant the kinematic terms $-\overline{L}\gamma^{\mu}\partial_{\mu}L-\overline{R}\gamma^{\mu}\partial_{\mu}R$ of the Lagrangian consists of the electronic isospin \overline{T} acting on L, plus the numbers N_L , N_R of left- and right-handed electron-type leptons. As far as we know, two of these symmetries are entirely unbroken: the charge $Q = T_3 - N_R - \frac{1}{2}N_L$, and the electron number $N=N_R+N_L$. But the gauge field corresponding to an unbroken symmetry will have zero mass,4 and there is no massless particle coupled to N,5 so we must form our gauge group out of the electronic isospin \vec{T} and the electronic hyperchange $Y = N_R$ $+\frac{1}{2}N_{L}$.

Therefore, we shall construct our Lagrangian out of L and R, plus gauge fields \hat{A}_{μ} and B_{μ} coupled to \vec{T} and Y, plus a spin-zero dou-

$$\varphi = \begin{pmatrix} \varphi^0 \\ \varphi^- \end{pmatrix}$$
 (3)

whose vacuum expectation value will break T and Y and give the electron its mass. The on-ly renormalizable Lagrangian which is invariant under T and Y gauge transformations is

$$\mathfrak{L} = -\frac{1}{4}(\partial_{\mu}\vec{\mathbf{A}}_{\nu} - \partial_{\nu}\vec{\mathbf{A}}_{\mu} + g\vec{\mathbf{A}}_{\mu} \times \vec{\mathbf{A}}_{\nu})^{2} - \frac{1}{4}(\partial_{\mu}B_{\nu} - \partial_{\nu}B_{\mu})^{2} - \overline{R}\gamma^{\mu}(\partial_{\mu} - ig'B_{\mu})R - L\gamma^{\mu}(\partial_{\mu}ig\vec{\mathbf{t}} \cdot \vec{\mathbf{A}}_{\mu} - i\frac{1}{2}g'B_{\mu})L$$

$$-\frac{1}{2} \stackrel{1}{1} \stackrel{1}{\partial}_{\mu} \varphi - ig \overline{A}_{\mu} \cdot \stackrel{7}{1} \varphi + i \stackrel{1}{2} g' B_{\mu} \varphi \stackrel{1}{1}^2 - G_{e} (\overline{L} \varphi R + \overline{R} \varphi^{\dagger} L) - M_{1}^{2} \varphi^{\dagger} \varphi + k (\varphi^{\dagger} \varphi)^{2}. \quad (4)$$

We have chosen the phase of the R field to make $G_{\mathcal{C}}$ real, and can also adjust the phase of the L and Q fields to make the vacuum expectation value $\lambda = \langle \phi^{\alpha} \rangle$ real. The "physical" φ fields are then φ^-

PHYSICAL REVIEW LETTERS

o vacuum expec-

and therefore the

d φ have no phys-

an is gauge invar-

ined isospin mation which where without change

ee that G_e is very tht be very large,

to replace φ ev-

 φ_* and φ^- have mass hat the Goldstone

We see immediately that the electron mass is λG_e . The charged spin-1 field is $(\varphi^0 - \varphi^{0\dagger})/i\sqrt{2}$. (5)

$$W_{\mu} = 2^{-1/2} (A_{\mu}^{1} + iA_{\mu}^{2})$$
 (8)

$$M_{W} = \frac{1}{2} \lambda g$$
. (9)

The neutral spin-1 fields of definite mass are

$$Z_{\mu} = (g^2 + g^2)^{-1/2} (gA_{\mu}^3 + g^2B_{\mu}),$$
 (10)

$$A_{u} = (g^{2} + g^{2})^{-1/2}(-g'A_{u}^{3} + gB_{u}). \tag{11}$$

$$M_Z = \frac{1}{2}\lambda(g^2 + g^{\prime 2})^{1/2},$$
 (12)

$$M_A = 0$$
, (13)

so A_{μ} is to be identified as the photon field. The interaction between leptons and spin-1

$$_{\mu}^{+}$$
 H.c. $+\frac{igg^{*}}{(g^{2}+g^{\prime 2})^{1/2}}\bar{e}^{*}\gamma^{\mu}eA_{\mu}$
 $+\frac{i(g^{2}+g^{\prime 2})^{1/2}}{4}\left[\left(\frac{3g^{\prime 2}-g^{2}}{2g^{\prime 2}+g^{2}}\right)\bar{e}^{*}\gamma^{\mu}e-\bar{e}^{*}\gamma^{\mu}\gamma_{5}e+\bar{\nu}\gamma^{\mu}(1+\gamma_{5})\nu\right]Z_{\mu}.$ (1)

$$=1/2\lambda^2$$
.

/2 = 2.07×10⁻⁶

y weak. Note alrger than e, so while (12) gives

by this model have to do with the couplings of the neutral intermediate meson Z_μ . If Z_μ does not couple to hadrons then the best place to look for effects of Z_{μ} is in electron-neutron scattering. Applying a Fierz transformation to the W-exchange terms, the total effective $e-\nu$ interaction is

If $g\gg e$ then $g\gg g'$, and this is just the usual $e-\nu$ scattering matrix element times an extra factor $\frac{1}{2}$. If $g \simeq e$ then $g \ll g'$, and the vector interaction is multiplied by a factor - 2 rather than 3/2. Of course our model has too many arbitrary features for these predictions to be

scussed by S. Glashow, Nucl. Phys. 22, 579 the chief difference is that Glashow introduc king terms into the Lagrangian, and re gets less definite predictions oldstone, Nuovo Cimento 19, 154 (1961); J. Gold-A. Salam, and S. Weinberg, Phys. Rev. 127,

Physik 88, 161 (1934). A model similar to ours

W. Higgs, Phys. Letters <u>12</u>, 132 (1964), Phys. etters <u>13</u>, 508 (1964), and Phys. Rev. <u>145</u>, 1156 ; F. Englert and R. Brout, Phys. Rev. Letters (1964); G. S. Guralnik, C. R. Hagen, and T. W. e, Phys. Rev. Letters 13, 585 (1964). articularly T. W. B. Kibble, Phys. Rev. 155, 67). A similar phenomenon occurs in the nteractions; the ρ-meson mass in zeroth-order tion theory is just the bare mass, while the picks up an extra contribution from the sponson picks up an extra contribution from the spos-se breaking of chiral symmetry. See S. Weinberg, Rev. Letters 18, 507 (1967), especially footnote Schwinger, Phys. Letters 24B, 473 (1967); show, H. Schmitzer, and S. Weinberg, Phys. Rev. is 19, 139 (1967), Eq. (13) et seq. D. Lee and C. N. Yang, Phys. Rev. 98, 101 (1965). Is is the same sort of transformation as that eliminates the nonderivative 7 couplings in the el; see S. Weinberg, Phys. Rev. Letters 18, 188. The 7 reappears with derivative coupling be-The # reappears with derivative coupling be-the strong-interaction Lagrangian is not invarier chiral gauge transforms a similar argument applied to the σ meson, see

P. Feynman and M. Gell-Mann, Phys. Rev. 109,

KING, AND LEPTON-PAIR

ing the leptonic decay rates of ρ^0 ,

ed to the (1+8) vector currents of the



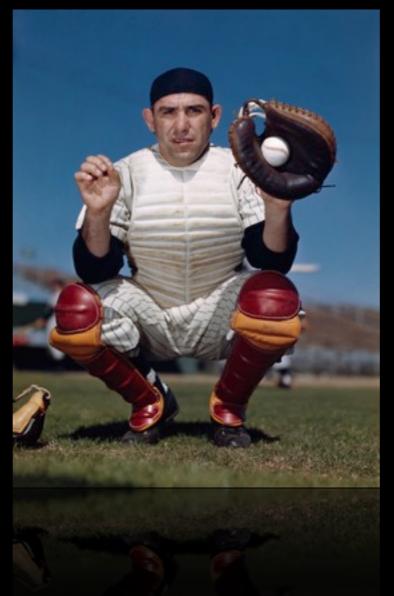
1967

Where Quality Starts Fresh Every Day









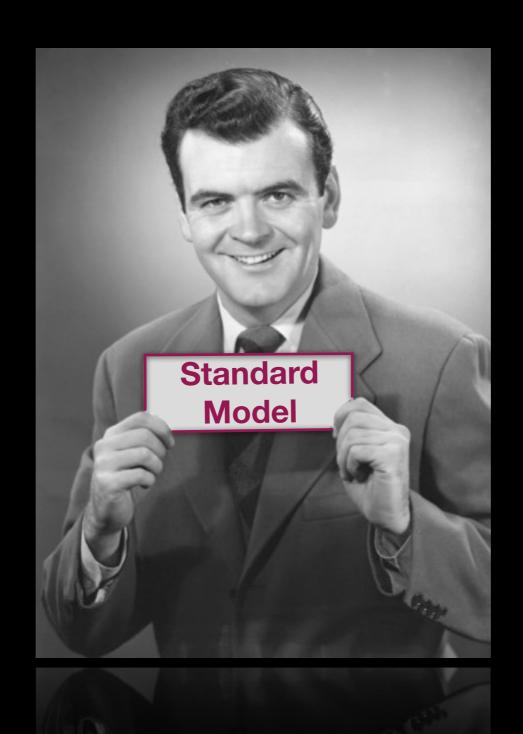
Ninety percent of this game is half mental.

The most precise theory in the history of physics

Quantity	Value	Standard Model	Pull	Dev.
M_Z [GeV]	91.1876 ± 0.0021	91.1874 ± 0.0021	0.1	0.0
Γ_Z [GeV]	2.4952 ± 0.0023	2.4961 ± 0.0010	-0.4	-0.2
Γ(had) [GeV]	1.7444 ± 0.0020	1.7426 ± 0.0010	_	_
$\Gamma(inv)$ [MeV]	499.0 ± 1.5	501.69 ± 0.06	_	_
$\Gamma(\ell^+\ell^-)$ [MeV]	83.984 ± 0.086	84.005 ± 0.015	_	_
$\sigma_{ m had} [m nb]$	41.541 ± 0.037	41.477 ± 0.009	1.7	1.7
R_e	20.804 ± 0.050	20.744 ± 0.011	1.2	1.3
R_{μ}	20.785 ± 0.033	20.744 ± 0.011	1.2	1.3
$R_{ au}$	20.764 ± 0.045	20.789 ± 0.011	-0.6	-0.5
R_b	0.21629 ± 0.00066	0.21576 ± 0.00004	0.8	0.8
R_c	0.1721 ± 0.0030	0.17227 ± 0.00004	-0.1	-0.1
$A_{FB}^{(0,e)}$	0.0145 ± 0.0025	0.03 ± 0.00021	-0.7	-0.7
$A_{FB}^{(0,\mu)}$	0.0169 ± 0.0013		0.4	0.6
$A_{FB}^{(0, au)}$	0.0188 ± 0.0017		1.5	1.6
$A_{FB}^{(0,b)}$	0.0992 ± 0.0016	0.10 \.0007	-2.6	-2.3
$A_{FB}^{(0,c)}$	0.0707 ± 0.0035	1005	-0.9	-0.0
$A_{FB}^{(0,s)}$	0.0976 ± 0.01	07		
$\bar{s}_{\ell}^2(A_{FB}^{(0,q)})$	0.2324 ± 0.00			
	$0.23200 \pm 0.$			
	0.2287 ± 0.0			
A_e	0.15138 ± 0.0			
	0.1544 ± 0.006			
	0.1498 ± 0.004		0.5	0.0
A_{μ}	0.142 ± 0.015		-0.4	-0.3
A_{τ}	0.136 ± 0.015		-0.8	-0.7
	0.1439 ± 0.0043		-0.8	-0.7
A_b	0.923 ± 0.020	0.9348 ± 0.0001	-0.6	-0.6
A_c	0.670 ± 0.027	0.6680 ± 0.0004	0.1	0.1
A_s	0.895 ± 0.091	0.9357 ± 0.0001	-0.4	- 0.4

Quantity	Value	Standard Model	Pull	Dev.
m_t [GeV]	173.4 ± 1.0	173.5 ± 1.0	-0.1	-0.3
M_W [GeV]	80.420 ± 0.031	80.381 ± 0.014	1.2	1.6
	80.376 ± 0.033		-0.2	0.2
$g_V^{ u e}$	-0.040 ± 0.015	-0.0398 ± 0.0003	0.0	0.0
$g_A^{ u e}$	-0.507 ± 0.014	-0.5064 ± 0.0001	0.0	0.0
$Q_W(e)$	-0.0403 ± 0.0053	-0.0474 ± 0.0005	1.3	1.3
$Q_W(Cs)$	-73.20 ± 0.35	-73.23 ± 0.02	0.1	0.1
$Q_W(Tl)$	-116.4 ± 3.6	-116.88 ± 0.03	0.1	0.1
τ_{τ} [fs]	291.13 ± 0.43	290.75 ± 2.51	0.1	0.1
$\frac{1}{2}(g_{\mu}-2-\frac{lpha}{\pi})$	$(4511.07\pm0.77)\times10^{-9}$	$(4508.70\pm0.09)\times10^{-9}$	3.0	3.0

J. Beringer et al.(PDG), PR D86, 010001 (2012) (http://pdg.lbl.gov)



Standard

Model

```
standard |standard|
noun

1. a level of quality
or attainment
```

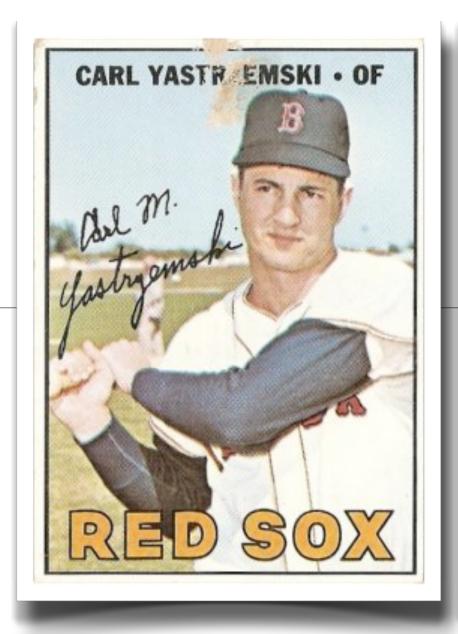
```
model | madl | noun | 2. a simplified | description, esp. a | mathematical one, of a | system or process, to | assist calculations and predictions
```

1967 - 2012

history was made

1967 - 2012

history was made







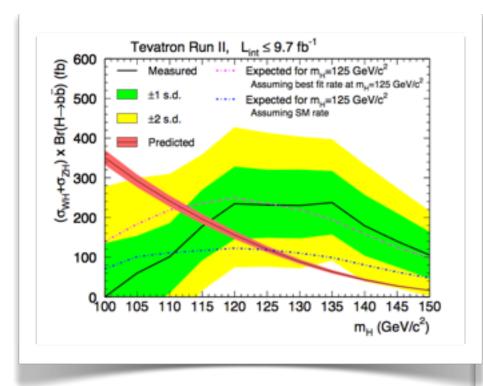


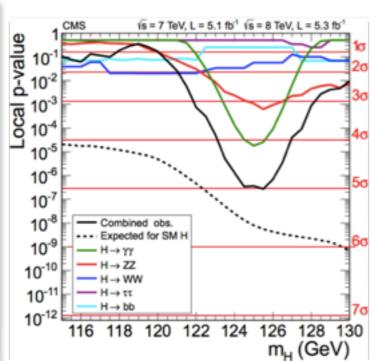


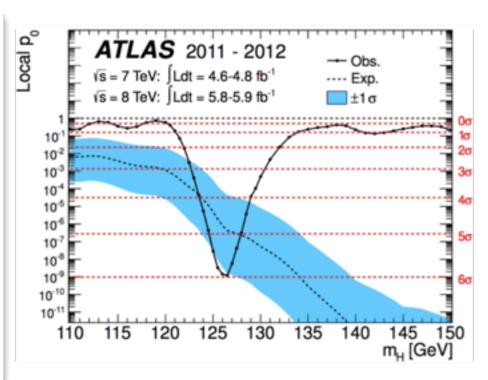






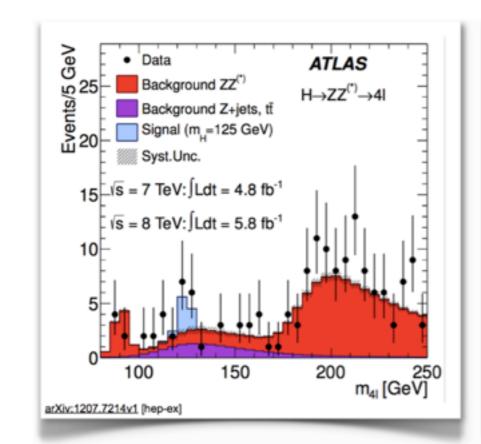


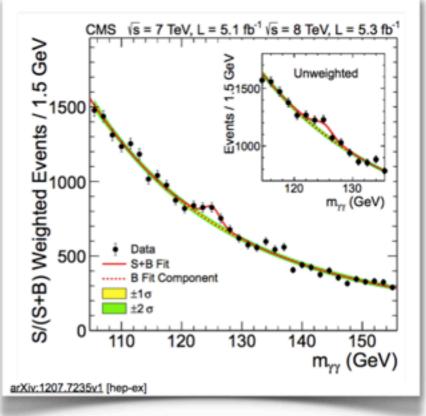




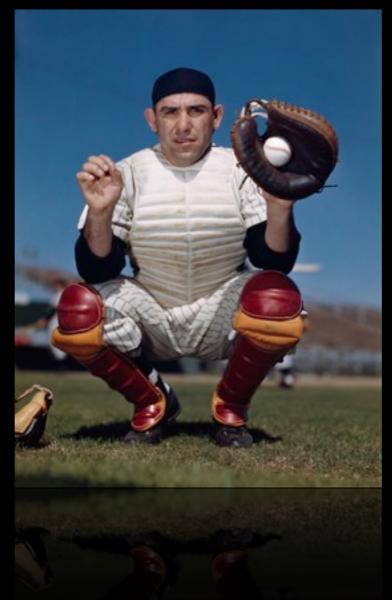
1967 - 2012

history was made





Peskin/Brock

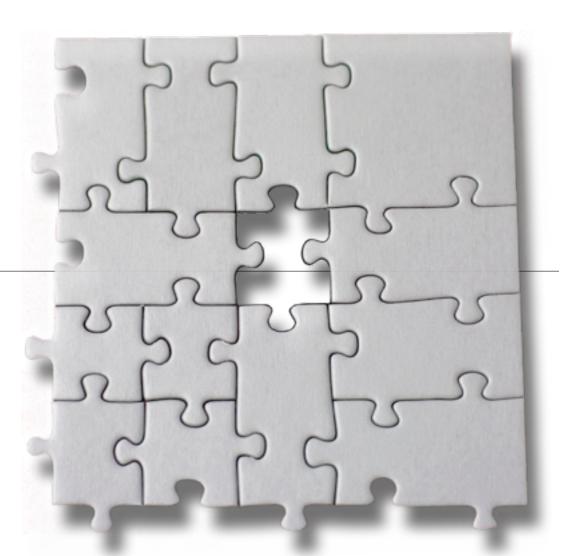


l always thought that record would stand until it was broken.



Standard Model

tempting to think that the puzzle is solved



But we know better

The High Energy Frontier has 4 jobs:

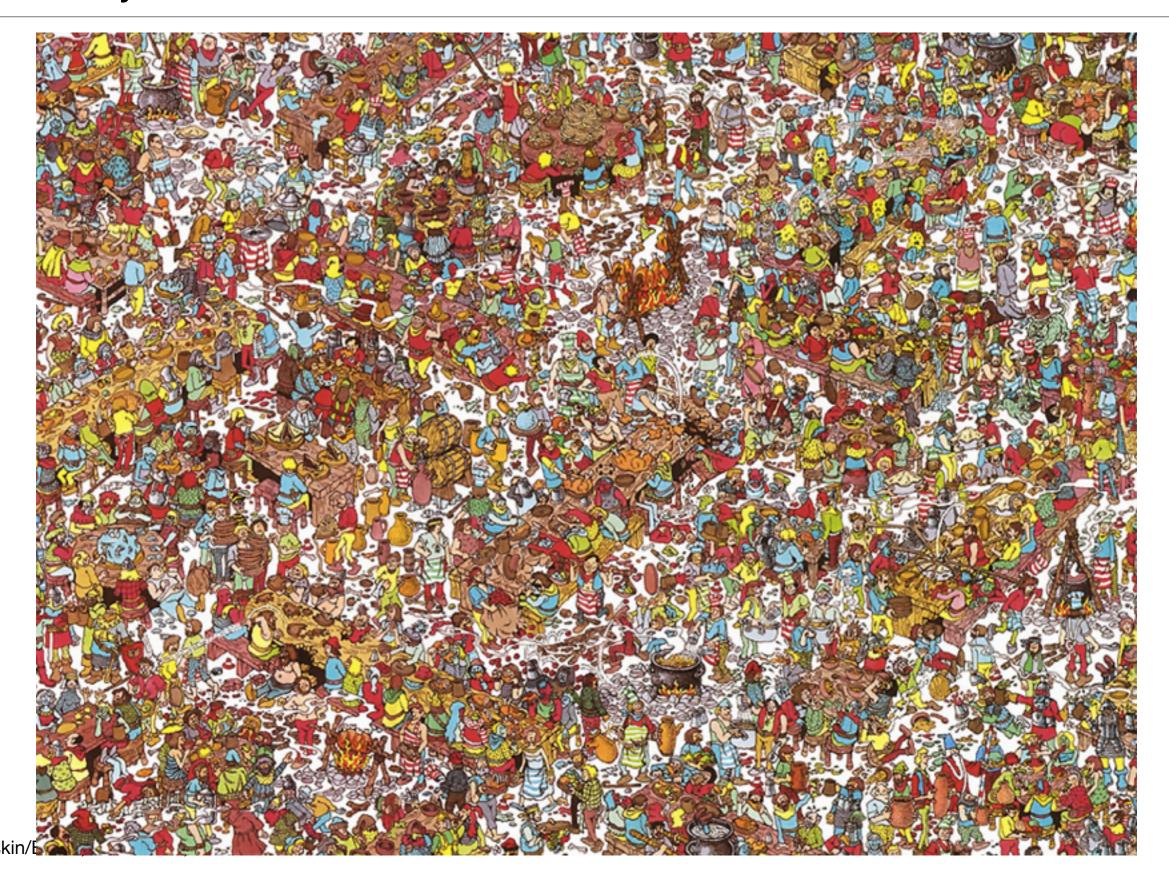
- Job 1. follow the obvious experimental imperative
- Job 2. unravel some confusions
- Job 3. complete a story
- Job 4. remember history

Job 1

what just happened?



the Object Itself





what's next with the Higgs-like-object?

HEF job: measure it... exhaustively



we have to measure:

mass

spin and parity (JP) looks like JP 0+,-?...2+,-? mixture?

couplings to vector bosons critical

all of the couplings to fermions, esp top and tau and b

couplings proportional to mass??

all branching fractions

singlet? how many are there?

elementary?

mixing with hidden sectors?

self-interactions?



Job 2

we have hair-on-fire observational

problems, disasters, conundrums, ...challenges.



First, the Higgs is too light



the same field theory

that in loops predicted the top mass range and now in loops the Higgs mass range

leads us to trouble opportunity

something has to protect the mass

quadratic divergences..."naturalness problem"

better: the "Naturalness Hint"

$$m_H - m_{\text{bare}} = \left(\frac{H}{H}\right) + \left(-\frac{t}{H}\right) + \left(-\frac{t}{H}\right) + \left(-\frac{t}{H}\right)$$

Taming this requires New Physics:

a symmetry? compositeness?

HEF job: unravel it

Second, galaxies move in ways that they shouldn't

Not like Newton

rotational velocity (km/s) measured

200

100

calculated

50000 100000 distance from center (light years)

Andromeda



since Nature's clumpy

there has to be a Dark Matter quantum





Third, where is all of the antimatter?



Third, where is all of the antimatter?

Some new force?

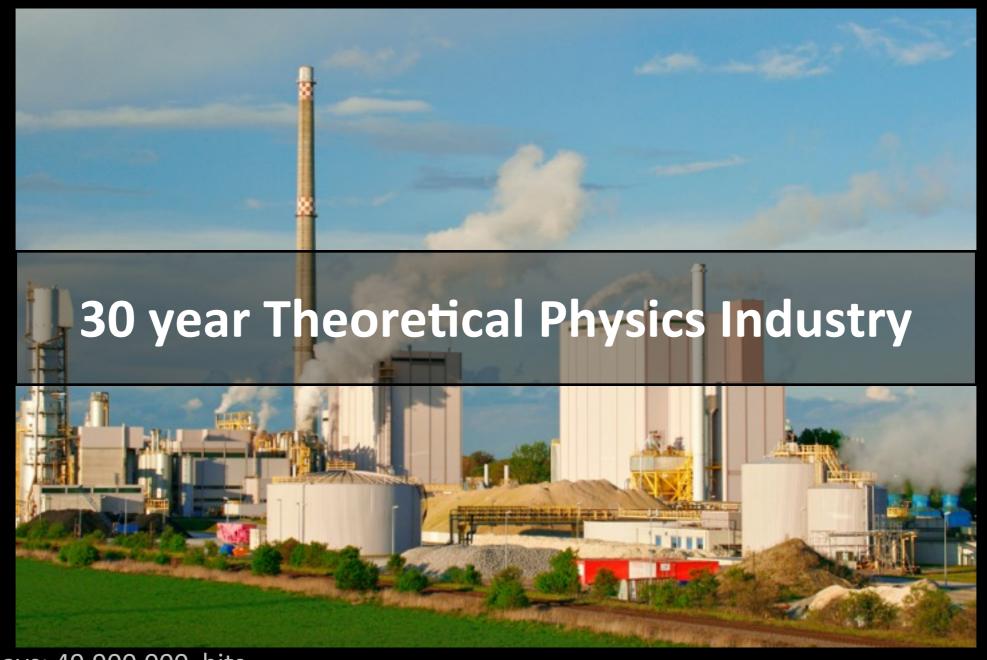
HEF job: find evidence





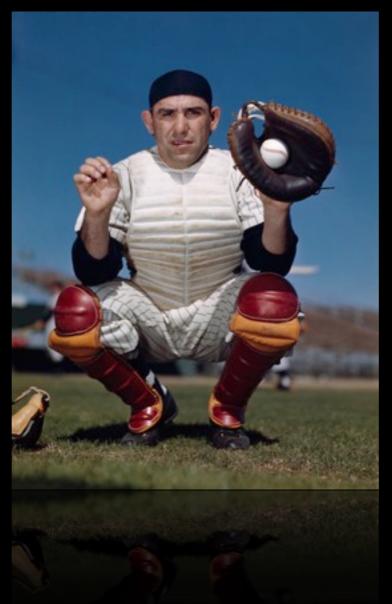


"Beyond the Standard Model"*



^{*} Mr Google says: 49,900,000 hits.





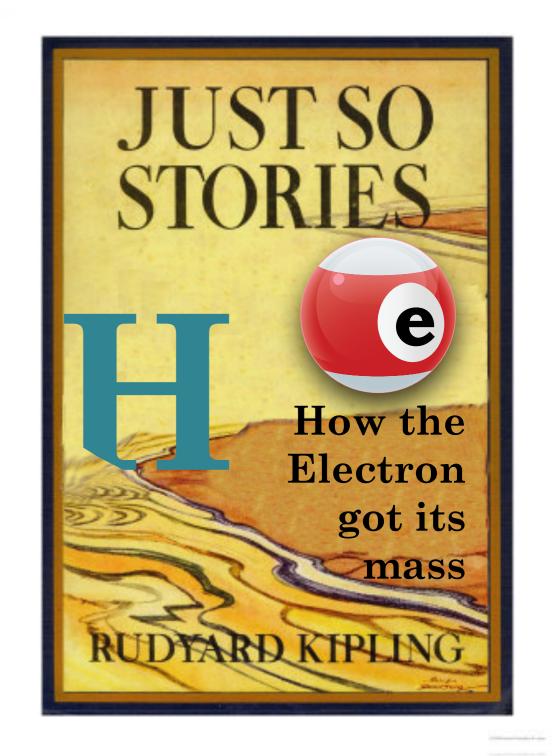
I never said most of the things I said.

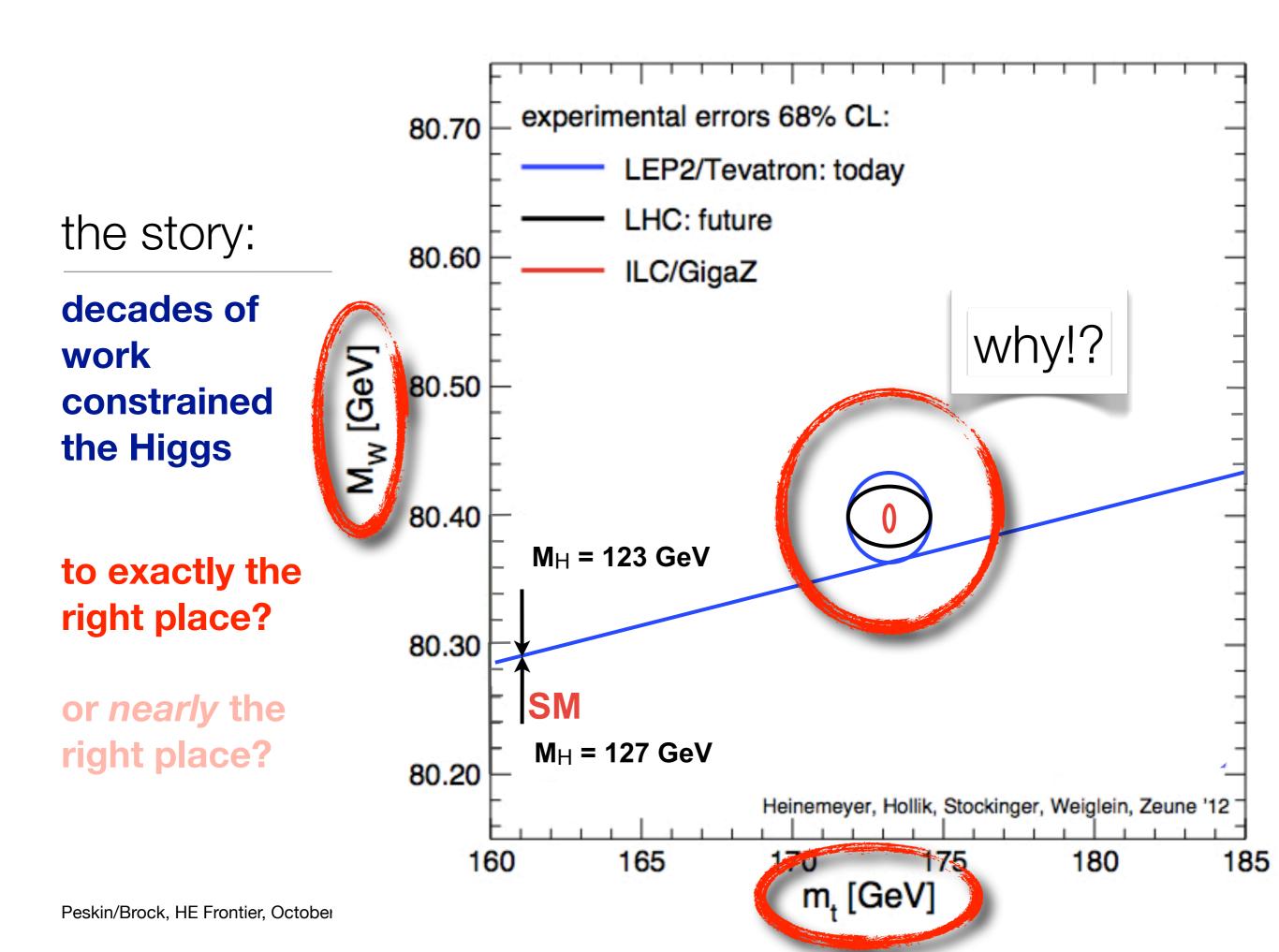
Job 3

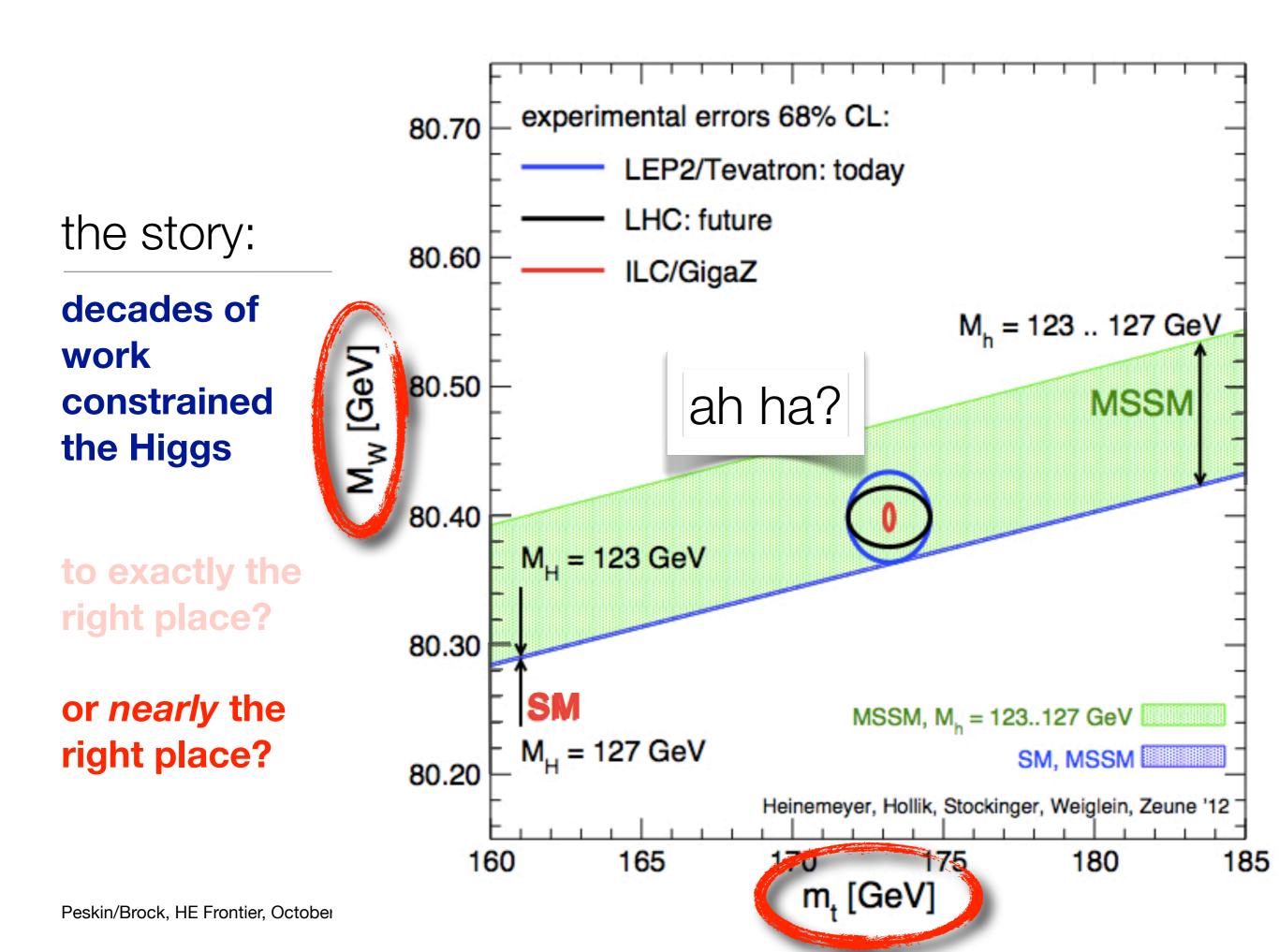
we have a story to tell

The Standard Model didn't have to unfold so...nicely

We don't know the characters, the plot is still developing, but the storyline is gripping.







SM requires correlations among measurables

inconsistencies have to show up



Job 4

things don't always follow a roadmap!

remember the:

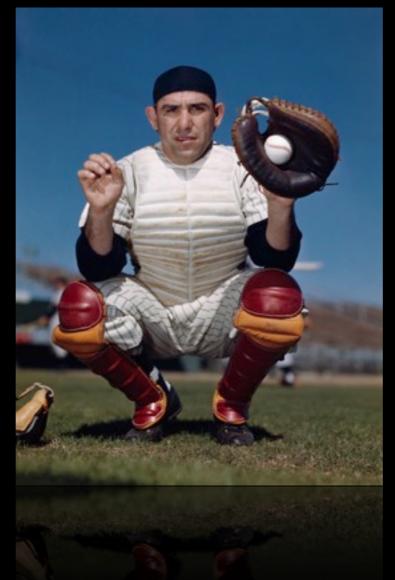
B lifetime

muon

Charm

etc



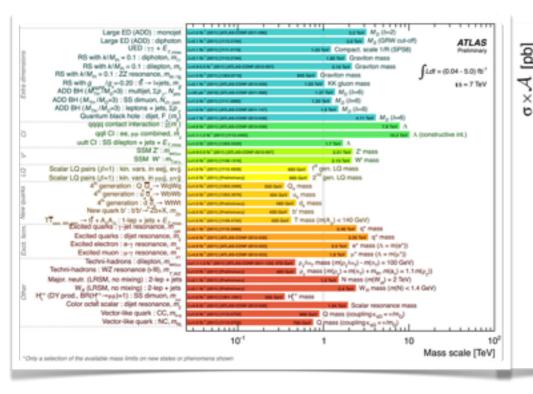


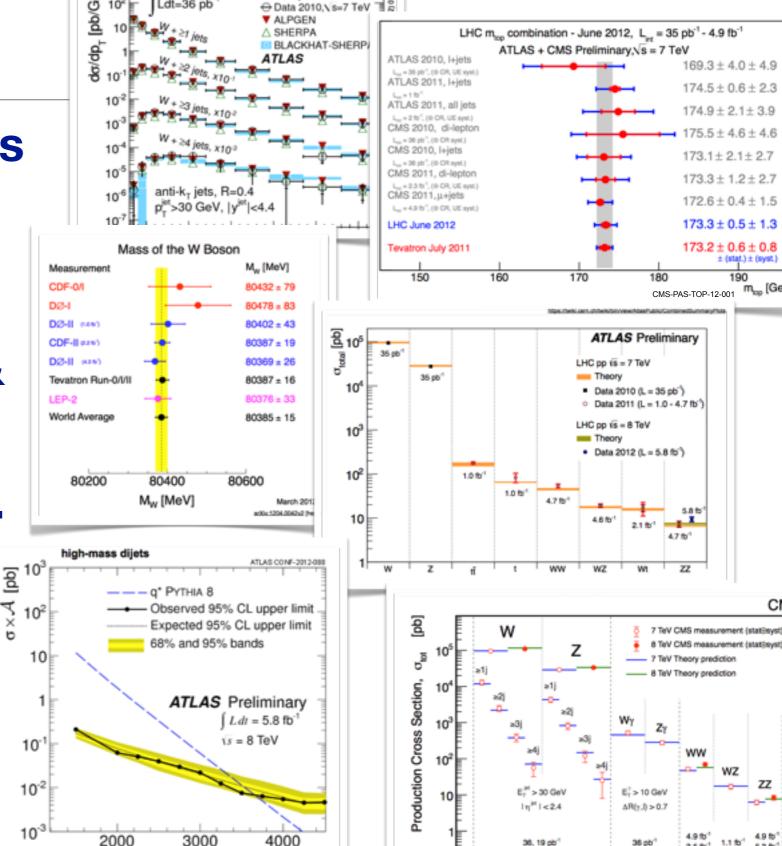
You've got to be very careful if you don't know where you are going because you might not get there.

Precision is the tool

Understanding of Jets Understanding Top W mass

Model independent & dependent searches and more...and more.





W→h' + jets

Mass [GeV]

1.1 fb⁻¹

36, 19 pb

CMS PKS SMP 12-011 DKZ 8 TeV

ZΖ

m_{top} [GeV]

CMS

W→h' + jets Data 2010, \s=7 TeV ALPGEN △ SHERPA LHC m_{son} combination - June 2012, L_{sc} = 35 pb⁻¹ - 4.9 fb⁻¹ Precision BLACKHAT-SHERP ATLAS + CMS Preliminary, s = 7 TeV ATLAS 2010, I+jets 169.3 ± 4.0 ± 4.9 ATLAS 2011, I+jets $174.5 \pm 0.6 \pm 2.3$ ATLAS 2011, all jets 174.9 ± 2.1 ± 3.9 L. + 2 fs*, (0-CR, UE evet.) **Understanding of Jets** CMS 2010. di-lepton $175.5 \pm 4.6 \pm 4.6$ 173.1 ± 2.1 ± 2.7 - - 75 rs* 10 CB syst 1 CMS 2011, di-lepton 173.3 ± 1.2 ± 2.7 - 2.3 to 1 (0 CR, UE evet.) CMS 2011, µ+jets 172.6 ± 0.4 ± 1.5 ">30 GeV, |yiot|<4.4 **Understanding Top** L., = 4.9 fb⁻¹, (0-CR, UE syst.) $173.3 \pm 0.5 \pm 1.3$ $173.2 \pm 0.6 \pm 0.8$ Mass of the W Boson M_w [MeV] W mass ATLAS Preliminary Model independent & The New Penysics is hiding in these p and more...and mare Observed 95% CL upper limit Expected 95% CL upper limit HEF job: don't be jaded! 10 10.2 Major, neutr. (LRSM, no mixing) : 2-lep + jet Mass [GeV

Mass scale [TeV]

So, the High Energy Frontier's Jobs:

1. Study the Higgs-like state at 125 GeV



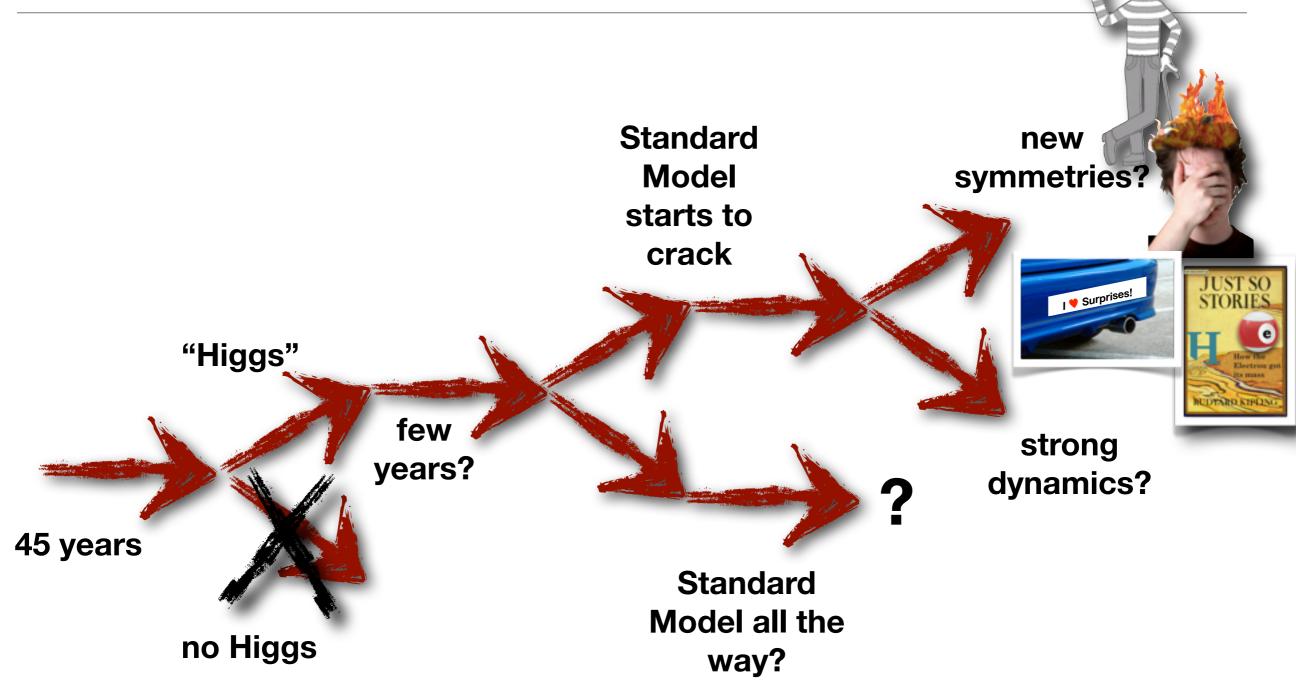
3. Write the story that encompasses the SM

4. Be nimble and observant to surprises





So, it's just getting interesting at the HEF!

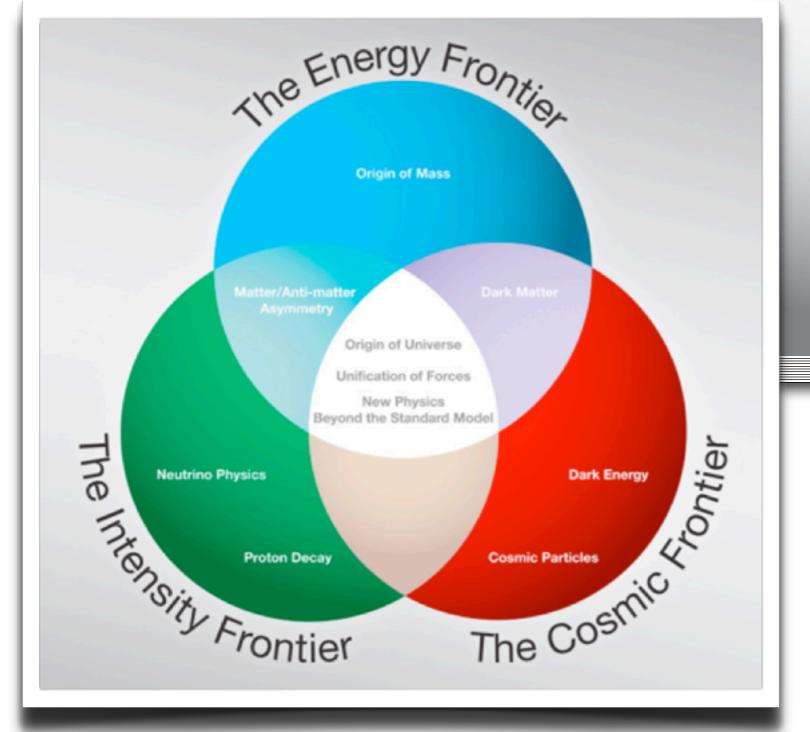


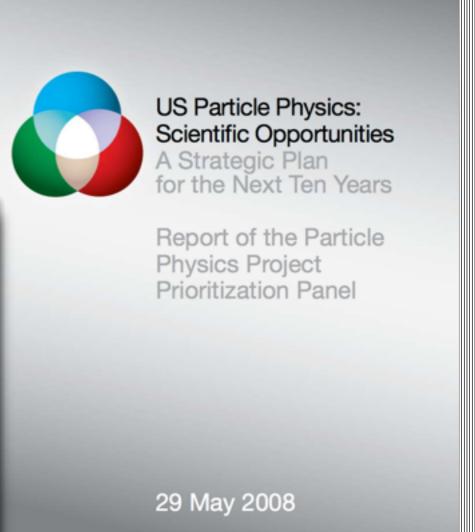
2. High Energy Frontier Study, practicalities

A gift from 2008 P5:

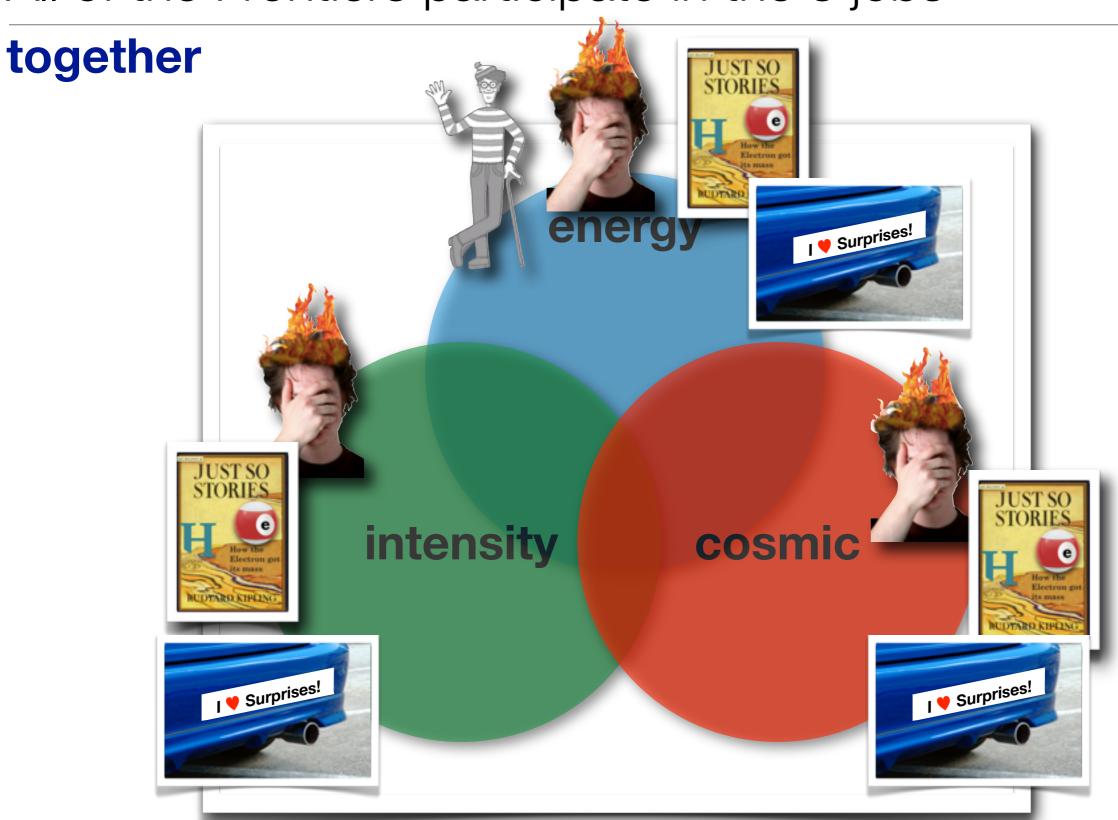
"the Frontier Circles"

and we're organized around them





All of the Frontiers participate in the 3 jobs



what we've done:

Identified terrific subgroup conveners

most have been meeting together for about a month

Created necessary correlations among groups

Decided on technical "connective tissue" groups

Explicit liaisons between HEF and other frontiers

Additional group "infrastructure"

established direct connection with the established collaborations:

"Contacts and consultants": ATLAS: Paul Tipton; CMS: Jim Olsen; LHCb: Sheldon Stone; ILD: Graham Wilson; SiD: Andy White; CLIC: Mark Thomson; Muon Collider: Ron Lipton

High Energy Frontier working groups

HE1: The Higgs Boson

Jianming Qian (Michigan), Andrei Gritsan (Johns Hopkins), Heather Logan (Carleton), Rick Van Kooten (Indiana), Chris Tully (Princeton), Sally Dawson (BNL)

HE2: Precision Study of Electroweak Interactions

Michael Schmitt (Northwestern), Doreen Wackeroth (Buffalo), Ashutosh Kotwal (Duke)

HE3: Fully Understanding the Top Quark

Robin Erbacher (Davis), Reinhard Schwienhorst (MSU), Kirill Melnikov (Johns Hopkins), Cecilia Gerber (UIC), Kaustubh Agashe (Maryland)

HE4: The Path Beyond the Standard Model–New Particles, Forces, and Dimensions

Daniel Whiteson (Irvine), Liantao Wang (Chicago), Yuri Gershtein (Rutgers), Meenakshi Narain (Brown), Markus Luty (UC Davis)

HE5: Quantum Chromodynamics and the Strong Interactions

Ken Hatakeyama (Baylor), John Campbell (FNAL), Frank Petriello (Northwestern), Joey Huston (MSU)

HE6: Flavor Physics and CP Violation at High Energy

Soeren Prell (ISU), Michele Papucci (LBNL), Marina Artuso (Syracuse)

HEF Goals:

1. In light of circa 2013 results what physics can be achieved before ~2018

...at design specifications with $\int \mathcal{L} dt \sim 100 \text{ fb}^{-1}$?

2. What are the LHC high luminosity physics goals for

... "Phase 1": circa 2022 with $\int \mathcal{L} dt$ of approximately 400 fb⁻¹

... "Phase 2": circa 2030 with $\int \mathcal{L} dt$ of approximately 3000 fb⁻¹

How do the envisioned upgrade paths inform those goals? Specifically, to what extent is precision Higgs Boson physics possible?

- 3. Does a Higgs Boson @ ~125 GeV/c² call for a "Higgs Factory"?
- 4. What are the physics cases for accelerators beyond 2025?

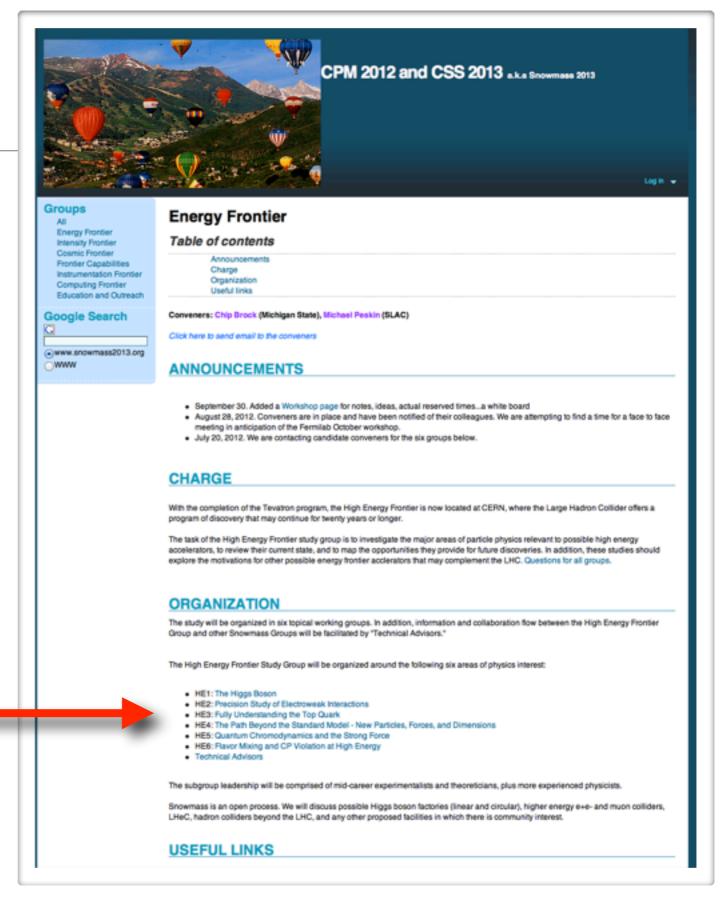
High energy LHC? High energy lepton collider? Lepton-hadron collider? VLHC?

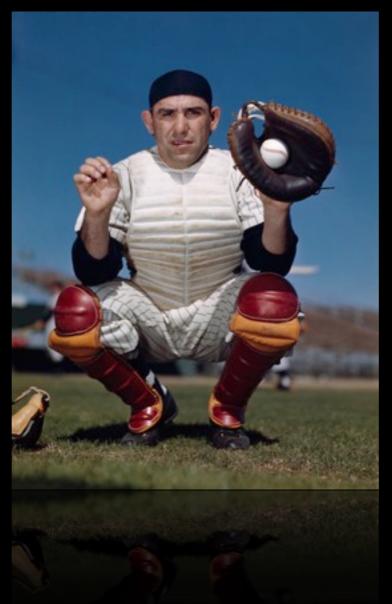
We too have a wiki

http://www.snowmass2013.org

increasingly active

find each group's detailed charges





Gyou can observe a lot just by watching.

Candidate scenarios to be addressed by all groups:

- A. The LHC with E = 14 TeV and $L = 10^{34}$ cm⁻² sec⁻¹ \checkmark
- B. A luminosity upgraded LHC with: $E_{cm} = 14$ TeV, $L = \sim 10^{35}$ cm⁻²s⁻¹ \checkmark
- C. An energy upgraded LHC
- D. e+e- lepton colliders $E_{cm} < \sim 1$ TeV \checkmark
- E. A circular e+e- collider operating as a Higgs factory.
- F. e+e- or gamma-gamma collider $E_{cm} > \sim 1$ TeV \checkmark
- G. A mu+mu- collider.
- H. A lepton-hadron collider. 🗸
- I. A VLHC hadron collider with energy well above the LHC energy. 🗸
- It is important to point out critical points in energy or luminosity that are essential to realize physics goals.
- For experiments at hadron colliders, a specific question is the effect of the machine environment for highluminosity running. Do high-luminosity conditions compromise the needed measurements? Are there detector designs or experimental strategies that can ameliorate these problems?

Common template Charge to each HEF Group:

1.Please provide a compact summary of the state of the search for X physics, including information from LEP, the Tevatron, and the LHC.

2.Please address the following goals for X physics in the future:

...tailored list of questions/goals follow

3.Please guide your exploration of the above goals with the following scenarios/caveats:

- Evaluate the above goals in the context of Candidate Facilities A-I. (Collaboration with the Facilities Frontier is expected.)
- Are new theoretical or simulation tools (for signal or backgrounds) required in order to achieve the goals?
- What are the detector and computing challenges that the above goals imply? (Collaboration with the Instrumentation and Computing Frontiers is expected.)

an expurgated version

of their charges

The Higgs Boson:

- How will we measure the full phenomenological profile of the Higgs boson?
- What level of precision can be achieved at the various proposed accelerators?
- What are the unique capabilities of each program?
- How will we discover possible additional states in the Higgs sector?
- To what extent are properties of the Higgs sector important more generally for fundamental physics?

Precision Study of Electroweak Interactions:

 What are the most important precision observables that will be studied at proposed accelerators?

 What level of precision can be achieved, and what is the importance of these measurements?

 How well can we probe the couplings of the W and Z bosons?

 What do we hope to learn from these measurements?

Fully Understanding the Top Quark:

- How well can we measure the top quark mass and width at proposed accelerators?
- How well can we measure the couplings of the top quark?
- How deeply can we probe for rare decays of the top quark?
- How can we use these measurements to search for new physics?
- Are there new particles that decay to top? How can we find them?

The Path Beyond the Standard Model:

- What is the new picture of physics at the TeV scale including the new information from LHC?
- Can electroweak symmetry breaking still be "natural"? What does this imply?
- What types of new particles might be found at the various proposed accelerators?
- Are there more effective strategies to discover Supersymmetry, Composite Higgs, and other proposed models?
- How can accelerator experiments help to address the problem of dark matter?

Quantum Chromodynamics and the Strong Interactions:

 How can we improve the precision of our understanding of the strong interactions in perturbative QCD, in parton distributions, in nonperturbative physics?

How do we incorporate electroweak interactions into precision QCD?

 How can QCD concepts such as jet substructure be used as tools for experimental discovery?

Flavor Physics and CP Violation at High Energy:

 What are the viable models of TeV scale physics that include flavor non-universality and CP violation?

 What new particles or new signatures are implied by these theories? How will we discover them?

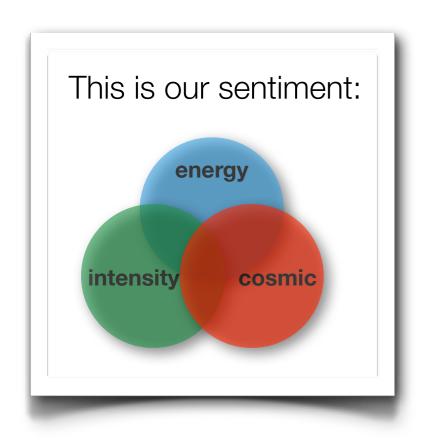
 How can high energy hadron colliders uniquely search for new physics in b and tau decays?

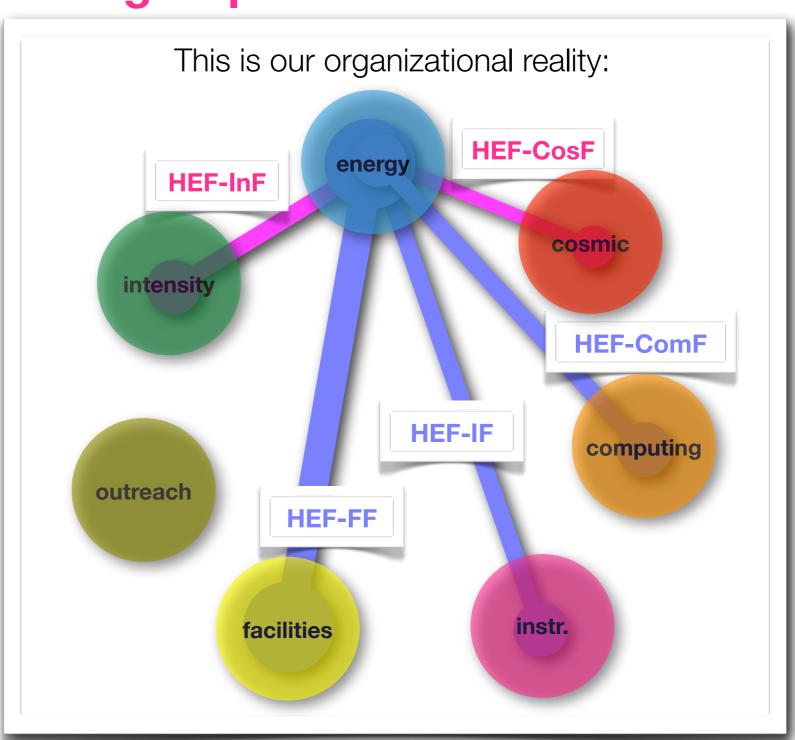
the overlaps

2 kinds of overlaps

Facilities, Instrumentation, and Computing Frontiers

Other Physics Frontiers groups





"technical group"

An explicit interface between the HEF physics groups and the FF, IF, and CF groups

Technical Group:

Beate Heinemann (Cal), Tom LeCompte (ANL), Jeff Berryhill (FNAL), Eric Torrence (Oregon), Tor Raubenheimer (SLAC), Eric Prebys (FNAL)

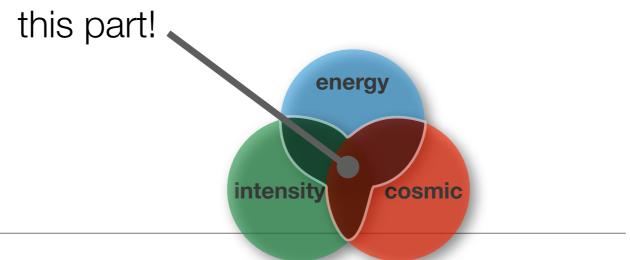
Early in the new year:

Establishing common benchmark parameters for each Candidate Facility

in support of the physics groups

Throughout the spring and workshop:

Liaison with the Facilities, Instrumentation and Computing Frontier Groups



Physics overlaps

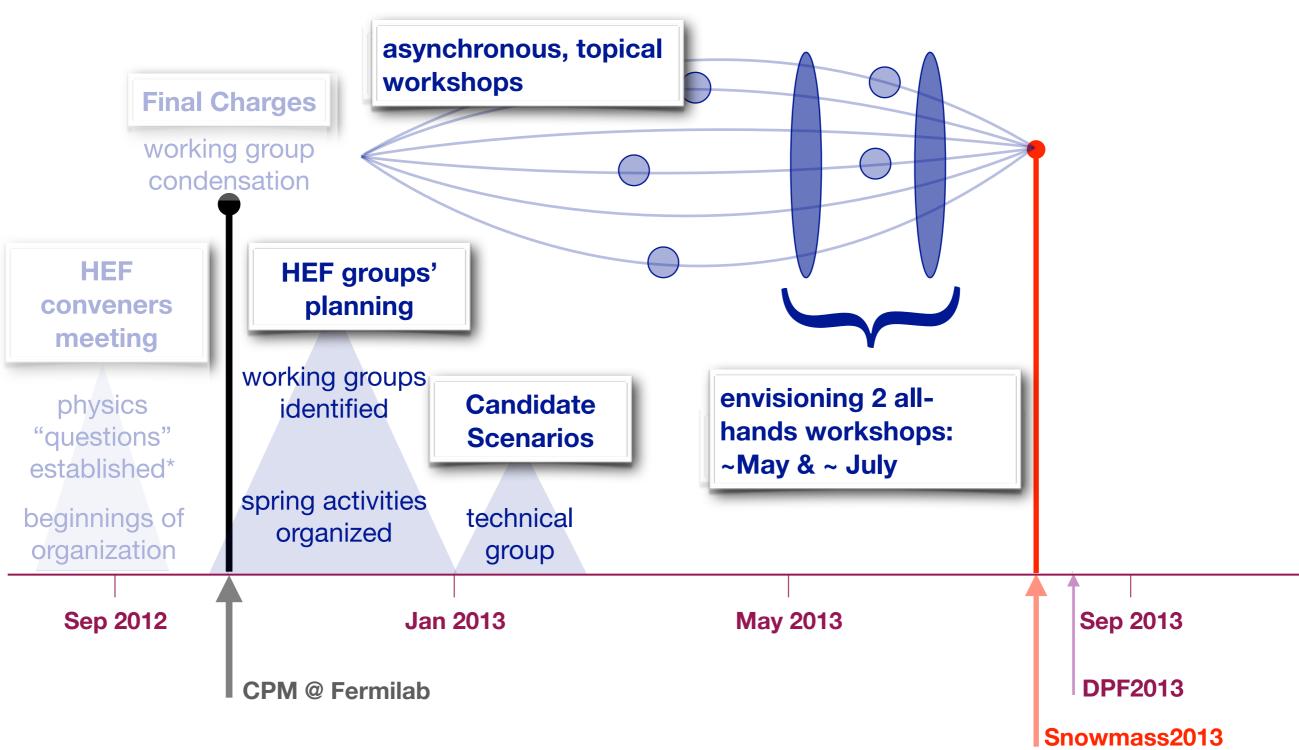
Explicit dual-coverage conveners:

HEF & CF (Dark Matter): Lian-Tao Wang & Konstantin Matchev HEF & CF (Baryogenesis): Michele Papucci & Ann Nelson HEF & HIF (b physics): Michele Papucci & Zoltan Ligeti

3. What's next for the High Energy Frontier

our to-do list

1. work.



4. This meeting: A full agenda for Friday



HEF Agenda for Friday, October 12th

During the day: The Auditorium

Subgroup Conveners will present

plans, organization, projects for study orientation to future colliders and challenges for experimentation

During the evening (19:30 - 21:00)

Six parallel working meetings, including some joint meetings ReadyTalk is available for all meetings.

During the Day: The Auditorium

Friday Morning

S1: 09:30 - 10:30 Landscape of Future Colliders

(joint HEF/Facilities)

Eric Prebys (FNAL): Hadron Colliders (15 + 5)

Mark Palmer (FNAL): Lepton Colliders (15+5)

Eric Torrence (Oregon): Physics Simulation for Hadron and Lepton Colliders (15+5)

S2: 10:45 - 11:45 Higgs / Electroweak

Chris Tully (Princeton): Higgs Physics at Hadron Colliders (15 + 5)

Rick Van Kooten (Indiana): Higgs Physics at Lepton Colliders (15 + 5)

Ashutosh Kotwal (Duke): Electroweak Physics after the Higgs Discovery (15+5)

S3: 12:00 - 13:00 Top / QCD

Joey Huston (Michigan State): Major Issues in QCD (25+5)

Reinhard Schwienhorst (Michigan State): Top quark physics (25+5)

(no break for lunch!)

During the Day: The Auditorium

Friday Afternoon

S4: 13:15 - 14:15 Challenges for Future Experiments

(joint HEF/Instrumentation/Computing)

Ulrich Heintz (Ohio): Detector Challenges for Future Hadron Colliders (15+5)

Andy White (UT Arlington): Detector Challenges for Future Lepton Colliders (15+5)

Ian Fisk (FNAL): Computing Challenges and Opportunites in the Next Decade (15+5)

S5: 14:30 - 15:30 **New Particles / Flavor**

Yuri Gershtein (Rutgers): What are the options for search strategies; how do we evaluate them? (5)

George Redlinger (BNL): Future of the search for supersymmetry at colliders (15)

Paddy Fox (FNAL): Future of the search for composite Higgs models and related exotic particles at colliders (15)

Michele Papucci (LBNL): Flavor Physics and CP violation at high energy (25)

During the Evening: Parallel Sessions

Friday Evening

19:30 - 21:00 New Particles, Daniel Whiteson

WH3NE (THEORY)

Informal discussion of new particle searches: What studies should be done, how do we cooperate between groups, what should the final product be?



WH3NW (Conjectorium)

Informal discussion of Flavor Physics and CP violation at high energy: What are the important topics to study?



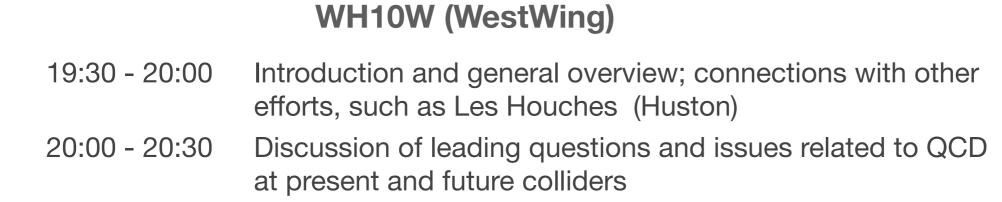


During the Evening: Parallel Sessions

Friday Evening

19:30 - 21:00	Top Quark, Reinhard Schwienhorst		
	WH7W (Racetrack)		
19:30 - 20:00	Discussion of precision top physics (leaders: Gerber + Melnikov)		
20:00 - 20:30	Discussion of detector physics involving top (leader: Schweinhorst)		
20:30 - 21:00	Discussion of new physics involving top (leaders: Erbacher and Agashe)		
19:30 - 21:00	QCD, Joey Huston		





Joint discussion with Electroweak group (leader: Petriello)



20:30 - 21:00

During the Evening: Parallel Sessions

Friday Evening

19:30 - 21:00 Higgs, Chris Tully
WH11NE (Sunrise)

19:30 - 21:00 Informal discussion of studies to be done on hadron and lepton collider strategies for Higgs studies

19:30 - 21:00 Electroweak, Michael Schmitt

WH11SE (RoundTable)

19:30 - 19:50 discussion of future electroweak precisions measurements (leader: Ashutosh Kotwal)
 19:50 - 20:10 Discussion of W and Z trilinear and Quartic couplings (leader: Michael Schmitt)
 20:10 - 20:30 Discussion of Electroweak fits including Higgs (leader: Doreen Wackeroth)

joint with QCD....go to WH10W





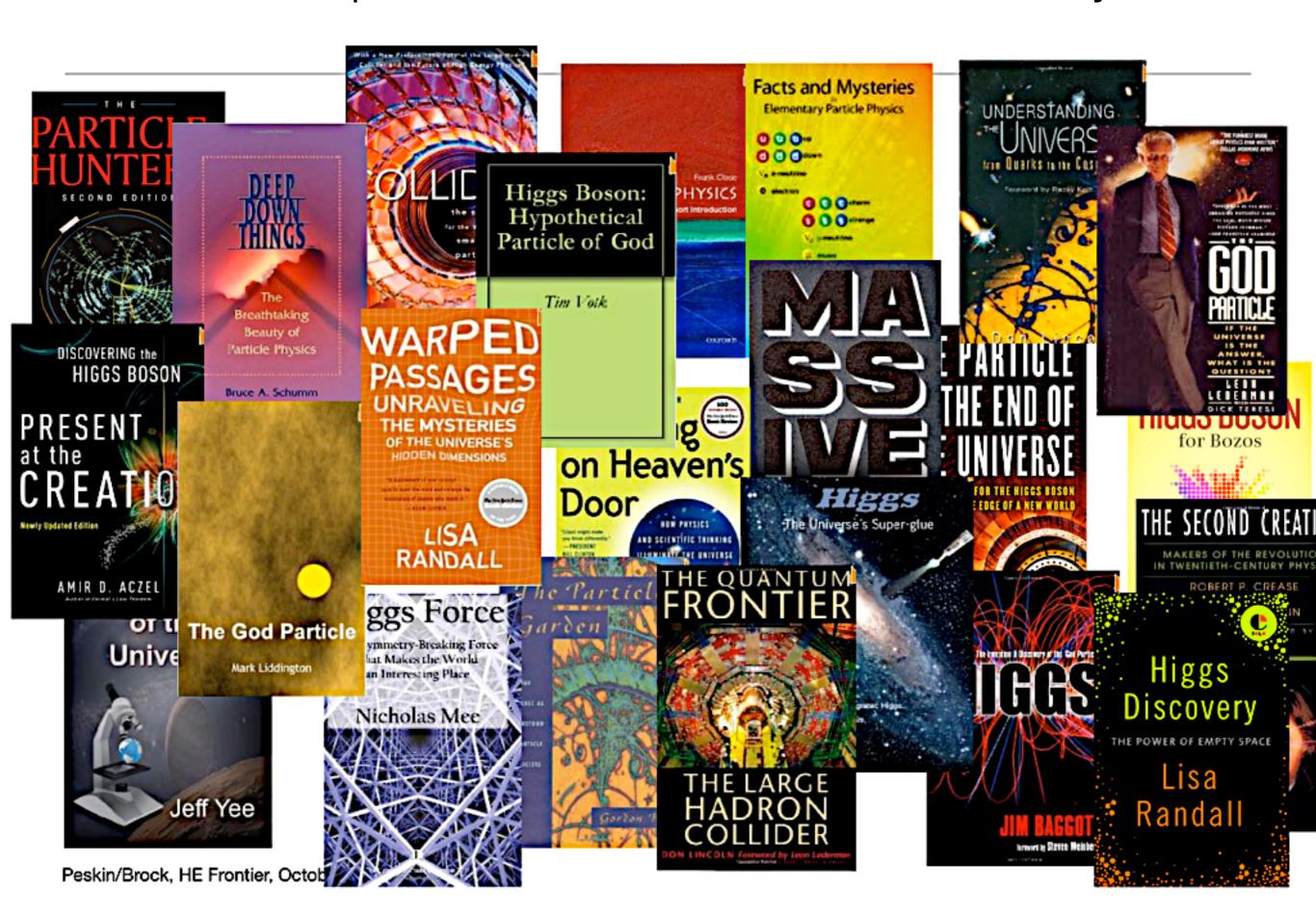
20:30 - 21:00

astonishingly

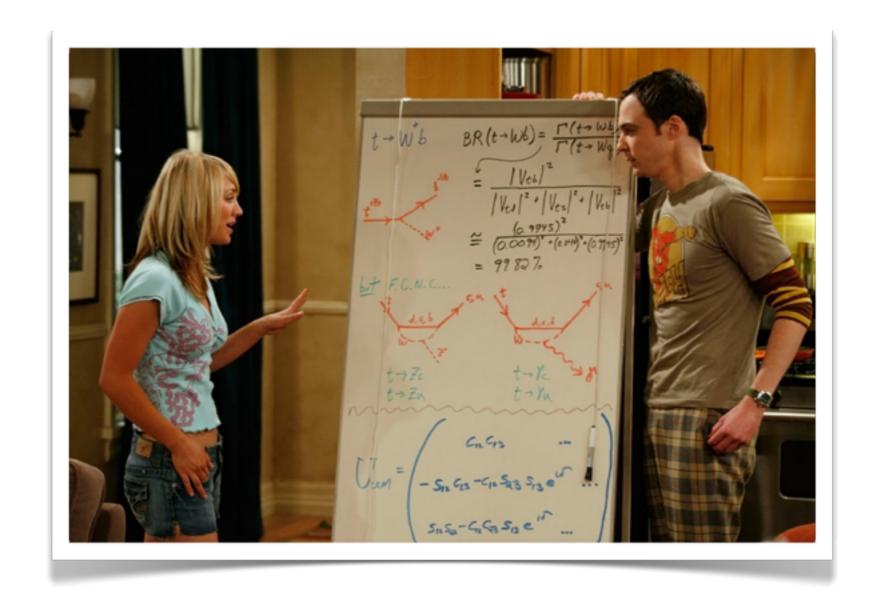
The world followed the Higgs Boson saga

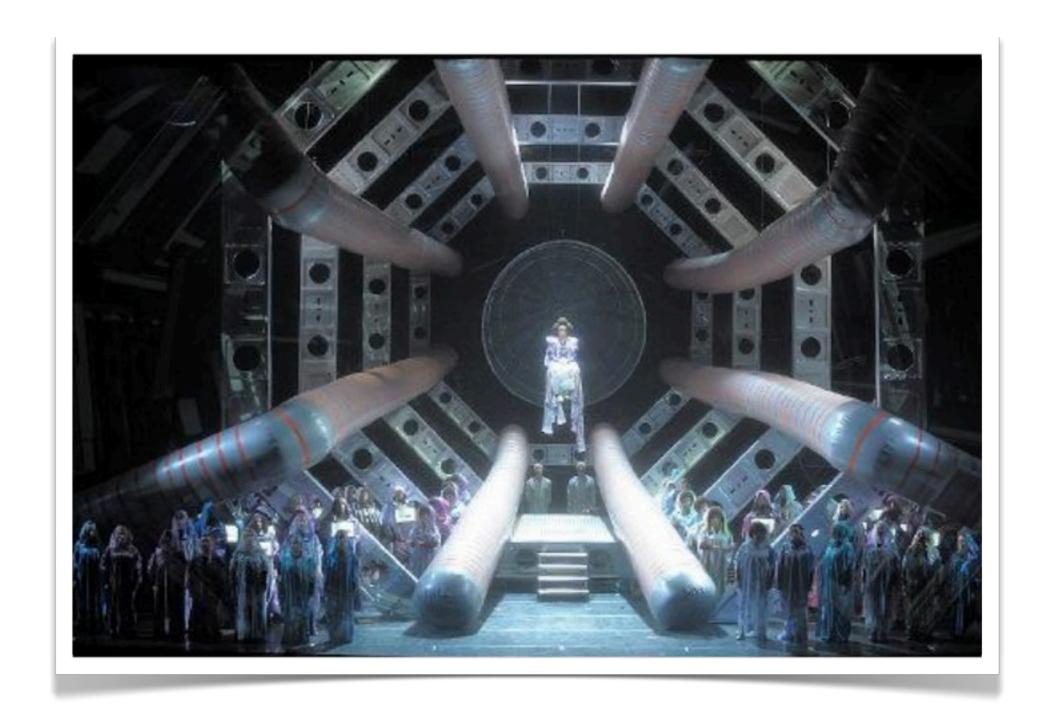
that's a big responsibility

a literature explosion that even astronomers envy

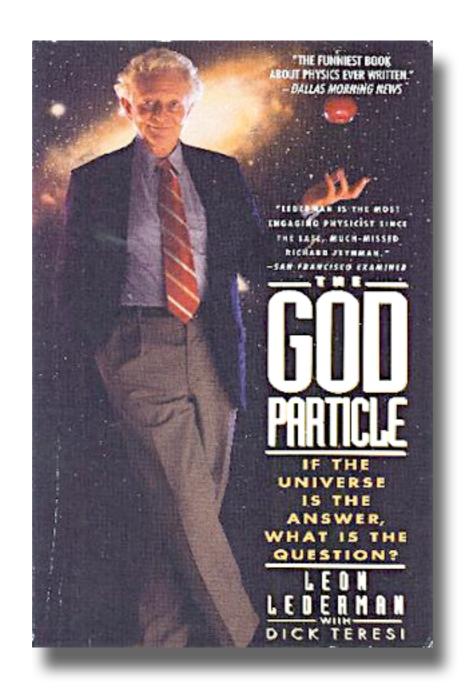














Because

the world has followed the Higgs Boson saga

we're in a much different place than in previous times.

Following the post-July physics is going to be fun.

Conclusions

Yes. This Snowmass is a big deal

U.S. participation is essential!

European and Asian participation is essential!

For HEF, the Higgs payoff will still be fresh

Snowmass2013 will guide HEPAP Strategic Planning Your colleagues need to be involved!



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